



Physics on $Y(5S)$ at B factories and a super B factory

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Outline

- Introduction.
- Recent Belle measurements at the Y(5S).
- Potential B_s studies at the Y(5S) at Belle and Super Belle.
- What else can be done at Super B factory?
- Conclusion.



Introduction

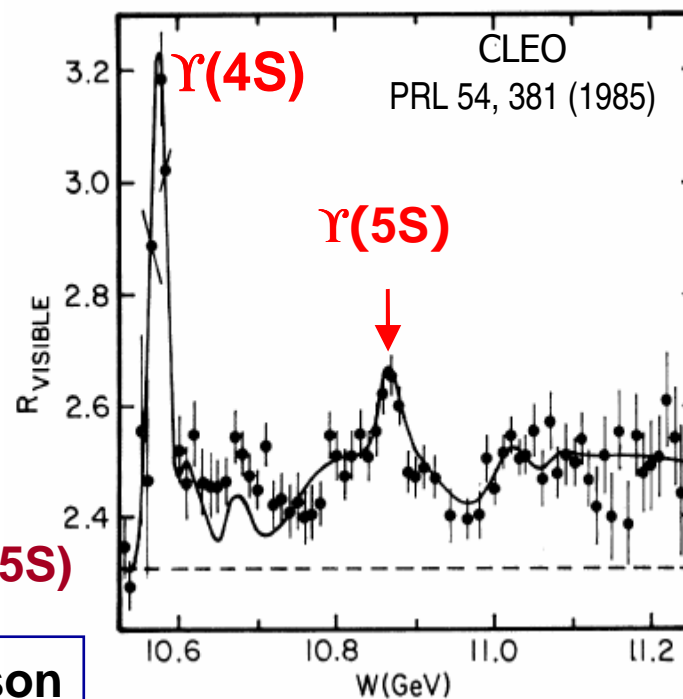
Asymmetric energy e^+e^- colliders
(**B Factories**) running at $Y(4S)$:
Belle and BaBar

1985: CESR (CLEO,CUSB) $\sim 116 \text{ pb}^{-1}$ at $Y(5S)$

2003: CESR (CLEO III) $\sim 0.42 \text{ fb}^{-1}$ at $Y(5S)$

2005: Belle, KEKB $\sim 1.86 \text{ fb}^{-1}$ at $Y(5S)$

2006, June 9-31: Belle, KEKB $\sim 21.7 \text{ fb}^{-1}$ at $Y(5S)$



$e^+ e^- \rightarrow Y(4S) \rightarrow B\bar{B}$, where B is B^+ or B^0 meson

$e^+ e^- \rightarrow Y(5S) \rightarrow B\bar{B}, B^*\bar{B}, B^*\bar{B}^*, B\bar{B}\pi, B\bar{B}\pi\pi, B_s\bar{B}_s, B_s^*\bar{B}_s, B_s^*\bar{B}_s^*$

where $B^* \rightarrow B \gamma$ and $B_s^* \rightarrow B_s \gamma$

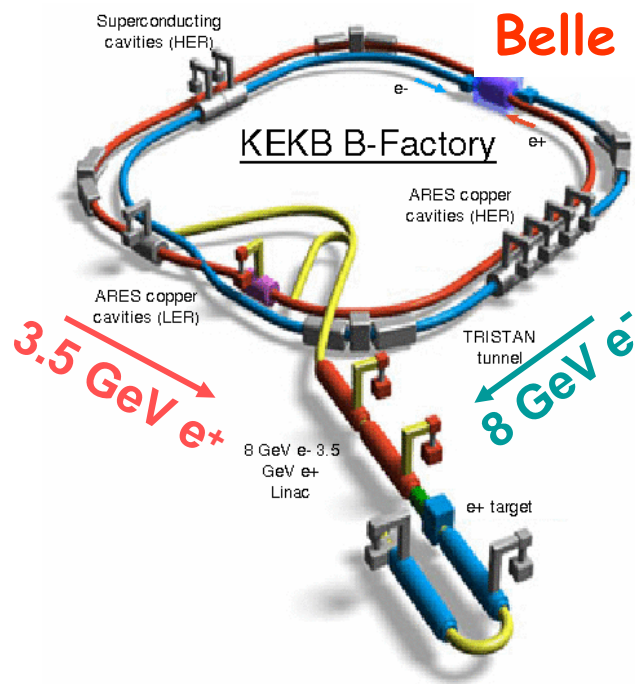
$M(Y(5S)) = 10865 \pm 8 \text{ MeV}/c^2$ (PDG)

$\Gamma(Y(5S)) = 110 \pm 13 \text{ MeV}/c^2$ (PDG)

B_s rate is $\sim 10\text{-}20\%$ \Rightarrow high lumi e^+e^- collider at the $Y(5S) \rightarrow B_s$ factory.



First Y(5S) runs at the KEKB e⁺e⁻ collider



Electron and positron beam energies were increased by 2.7% (same Lorentz boost $\beta\gamma = 0.425$) to move from Y(4S) to Y(5S).

No modifications are required for Belle detector, trigger system or software to move from Y(4S) to Y(5S).

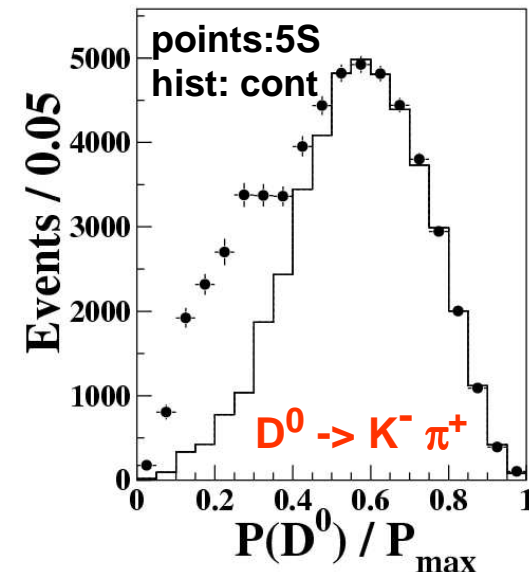
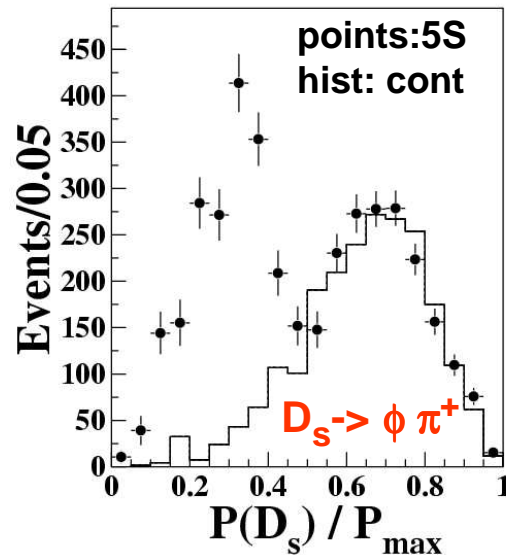
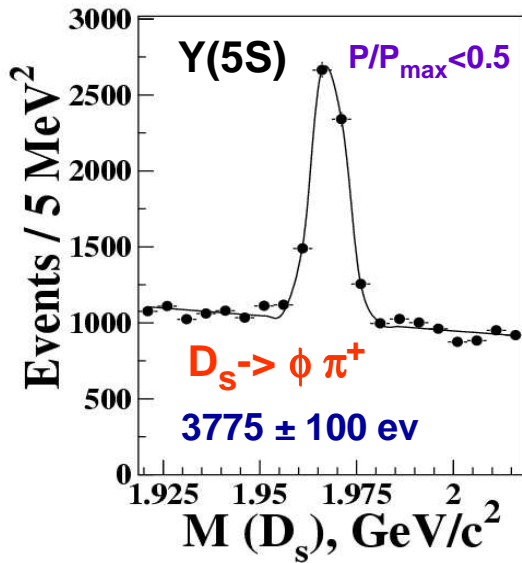
Integrated luminosity of $\sim 1.86 \text{ fb}^{-1}$ at 2005 and $\sim 21.6 \text{ fb}^{-1}$ at 2006 was taken by Belle detector at Y(5S).

The same luminosity per day can be taken at Y(4S) and Y(5S).

→ Very smooth running



Inclusive analyses : $Y(5S) \rightarrow D_s X$, $Y(5S) \rightarrow D^0 X$



After continuum subtraction and efficiency correction:

$$Bf(Y(5S) \rightarrow D_s X) / 2 = (23.6 \pm 1.2 \pm 3.6) \%$$

$$L = 1.86 \text{ fb}^{-1}$$

$$N_{bb}(5S) =$$

$$Bf(Y(5S) \rightarrow D^0 X) / 2 = (53.8 \pm 2.0 \pm 3.4) \%$$

$$561,000 \pm 3,000 \pm 29,000 \text{ events}$$

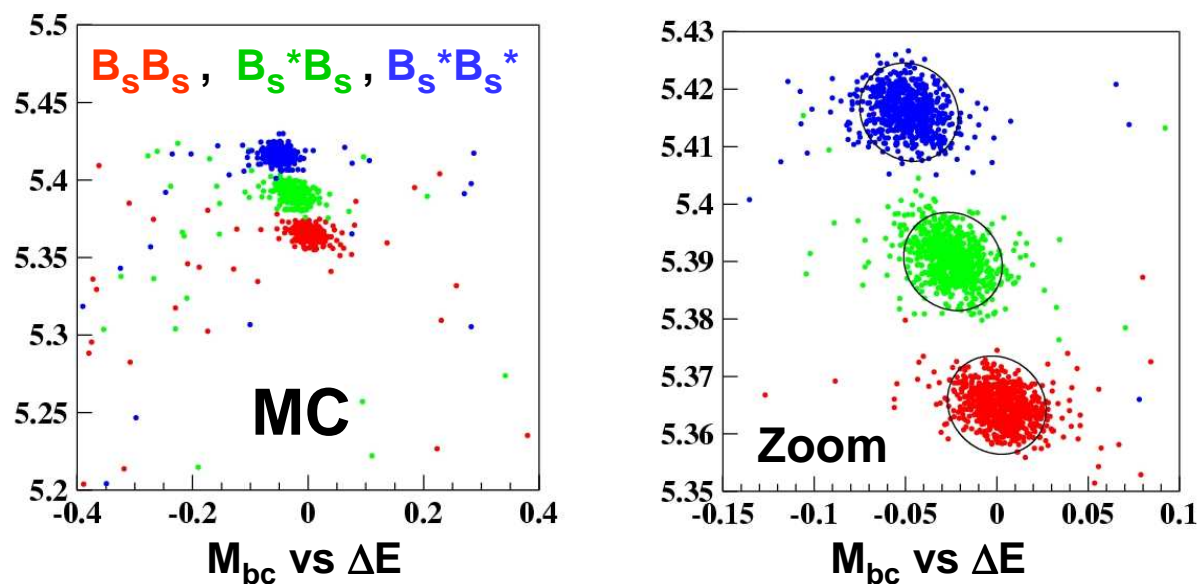
$$\Rightarrow f_s = N(B_s^{(*)} B_s^{(*)}) / N(bb) = (18.0 \pm 1.3 \pm 3.2) \%$$

$$N(B_s) / \text{fb}^{-1} = 108,000 \pm 21,000 \text{ events}$$

in good agreement with CLEO



Signature of fully reconstructed exclusive B_s decays



$e^+ e^- \rightarrow Y(5S) \rightarrow B_s B_s, B_s^* B_s, B_s^* B_s^*$, where $B_s^* \rightarrow B_s \gamma$

Reconstruction: B_s energy and momentum, photon from B_s^* is not reconstructed.

Two variables calculated: $M_{bc} = \sqrt{E_{beam}^{*2} - P_B^{*2}}$, $\Delta E = E_B^* - E_{beam}^*$

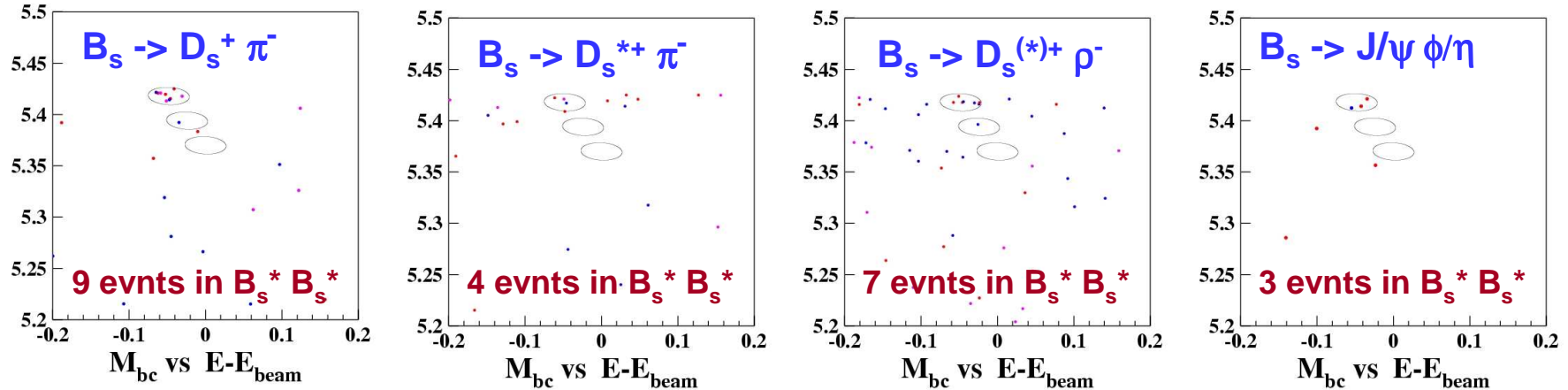
Figures (MC simulation) are shown for the decay mode $B_s \rightarrow D_s^- \pi^+$ with $D_s^- \rightarrow \phi \pi^-$.

The signals for $B_s B_s, B_s^* B_s$ and $B_s^* B_s^*$ can be well separated.



Exclusive $B_s \rightarrow D_s^{(*)+} \pi^- / \rho^-$ and $B_s \rightarrow J/\psi \phi / \eta$ decays

Data at $\Upsilon(5S)$, 1.86 fb^{-1}

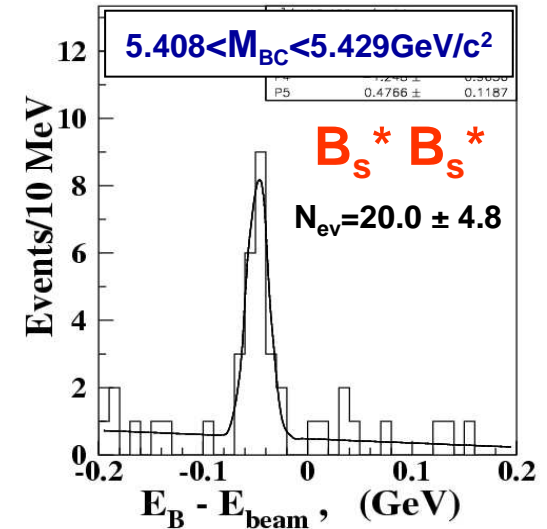


$$N(B_s^* B_s^*) / N(B_s^{(*)} B_s^{(*)}) = (94 \pm 6_9)\%$$

Potential models predict $B_s^* B_s^*$ dominance over $B_s^* B_s$ and $B_s B_s$ channels, but not so strong.

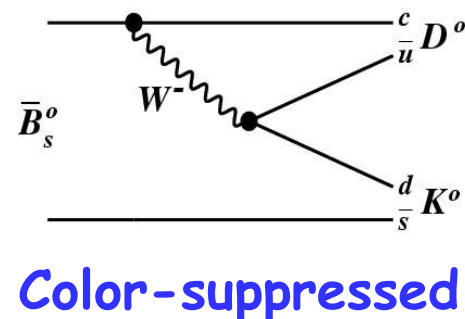
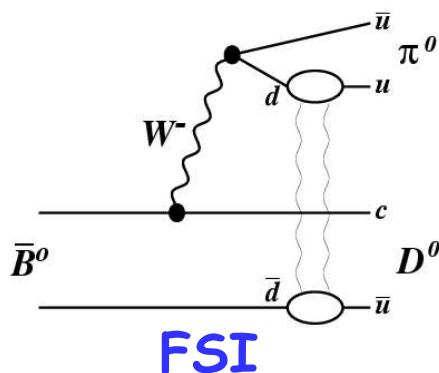
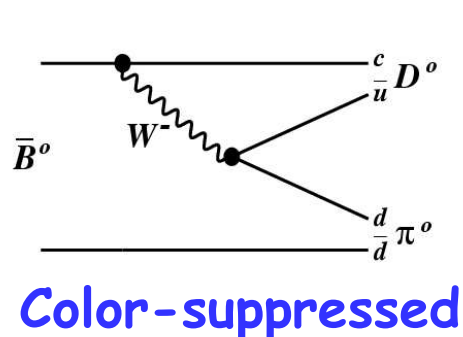
Conclusions:

1. Belle can take $\sim 30 \text{ fb}^{-1}$ per month (x2 soon).
2. Number of produced B_s at $\Upsilon(5S)$ is $\sim 10^5 / \text{fb}^{-1}$.
3. $B_s^* B_s^*$ channel dominates over all $B_s^{(*)} B_s^{(*)}$.
4. Backgrnds in exclusive modes are not large.





Color-suppressed $B_s \rightarrow D^0 K^0$ decay



$$\frac{Bf(B^0 \rightarrow D^0 \pi^0)}{Bf(B^0 \rightarrow D^+ \pi^-)} = \frac{(2.91 \pm 0.28) \times 10^{-4}}{(3.4 \pm 0.9) \times 10^{-3}} \approx 0.1$$

Which diagram, color-suppressed or FSI, is dominant in $B^0 \rightarrow D^0 \pi^0$ decay? Decay mode $B_s \rightarrow D^0 K^{(*)0}$ has no FSI diagram. If the ratio $Bf(B_s \rightarrow D^0 K^0) / Bf(B_s \rightarrow D_s^+ \pi^-) \sim 0.1$, then color-suppressed diagram dominates. If the ratio is significantly smaller, then FSI diagram dominates.

If $Bf(B_s \rightarrow D^0 K^0)$ is $\sim 3 \times 10^{-4}$, then ~ 8 events are expected with 25 fb^{-1} at $Y(5S)$.



Semileptonic B_s decays

At the $Y(5S)$ we can measure precisely semileptonic decays:

$$Bf(B_s \rightarrow X^+ \ell^- \nu)$$

$$Bf(B_s \rightarrow D_s^+ \ell^- \nu)$$

$$Bf(B_s \rightarrow D_s^{*+} \ell^- \nu)$$



Accuracy is expected to be
 $\sim(5-10)\%$ with 25 fb^{-1} at $Y(5S)$

Difficult to measure in hadron-hadron colliders.

These Bf s have to be compared with corresponding B meson Bf s. Within SM : $Bf(B_s \rightarrow X^+ \ell^- \nu) = Bf(B \rightarrow X^+ \ell^- \nu)$
If not, nonstandard contributions should be considered.

How to explain:

$\tau(B^0) > \tau(B_s)$ - 2.9σ difference (in contrast with theory).



$\Delta\Gamma_s/\Gamma_s$ measurement from $Bf (B_s \rightarrow D_s^{+(*)} D_s^{-(*)})$

$$M_{B_s} = (M_H + M_L)/2$$

$$\Gamma_s = (\Gamma_H + \Gamma_L)/2$$

$$\Delta m_s = M_H - M_L$$

$$\Delta\Gamma = \Gamma_L - \Gamma_H > 0 \text{ in SM}$$

$$i \frac{d}{dt} \begin{pmatrix} B_s \\ \bar{B}_s \end{pmatrix} = (M - i/2 \Gamma) \begin{pmatrix} B_s \\ \bar{B}_s \end{pmatrix} \quad - \text{ Schrödinger equation}$$

Matrices M and Γ are t-dependent, Hermitian 2x2 matrices

$$\text{Assuming CPT: } M_{11} = M_{22} \quad \Gamma_{11} = \Gamma_{22}$$

$$|B_{H,L}(t)\rangle = \exp(- (i M_{H,L} + \Gamma_{H,L}/2)t) |B_{H,L}\rangle$$

SM: $\beta_s = \arg(-V_{ts} V_{tb}^*/V_{cs} V_{cb}^*) = O(\lambda^2)$ - no CP-violation in mixing

$$\text{BSM: } \phi_s = \arg(-M_{12}/\Gamma_{12}) \quad 2\theta_s = \phi_s \quad \Delta\Gamma_s = 2 |\Gamma_{12}| \cos 2\theta_s$$



$\Delta\Gamma_s/\Gamma_s$ measurement from $Bf(B_s \rightarrow D_s^{+(*)} D_s^{-(*)})$

$$\Delta\Gamma_s = 2 |\Gamma_{12}| \cos \phi_s \quad \Delta\Gamma_s^{\text{SM}} = \Delta\Gamma_{\text{CP}^S} = 2 |\Gamma_{12}|$$

Since $\Delta\Gamma_{\text{CP}^S}$ is unaffected by NP, NP effects will decrease $\Delta\Gamma_s$.

$$\Delta\Gamma_{\text{CP}^S} = \sum \Gamma(\text{CP}=+) - \sum \Gamma(\text{CP}=-)$$

$B_s \rightarrow D_s^{(*)+} D_s^{(*)-}$ decays have CP- even final states with largest Bf 's of $\sim (1-3)\%$ each, saturating $\Delta\Gamma_s/\Gamma_s$.

$$\frac{\Delta\Gamma_{\text{CP}^S}}{\Gamma_s} \approx \frac{Bf(B_s \rightarrow D_s^{(*)+} D_s^{(*)-})}{1 - Bf(B_s \rightarrow D_s^{(*)+} D_s^{(*)-}) / 2}$$

To prove this formula experimentally : a) Contribution of $B_s \rightarrow D_s^{+(*)} D_s^{-(*)} n\pi$ is small b) Most of $B_s \rightarrow D_s^+ D_s^{-*}$ and $B_s \rightarrow D_s^{+*} D_s^-$ states are CP- even.

Assuming corrections are small ($\sim 5-7\%$), Bf measurement will provide information about $\Delta\Gamma_{\text{CP}^S}$ or $|\Gamma_{12}|$.



$\Delta\Gamma_s/\Gamma_s$ measurement from $Bf(B_s \rightarrow D_s^{(*)+} D_s^{(*)-})$

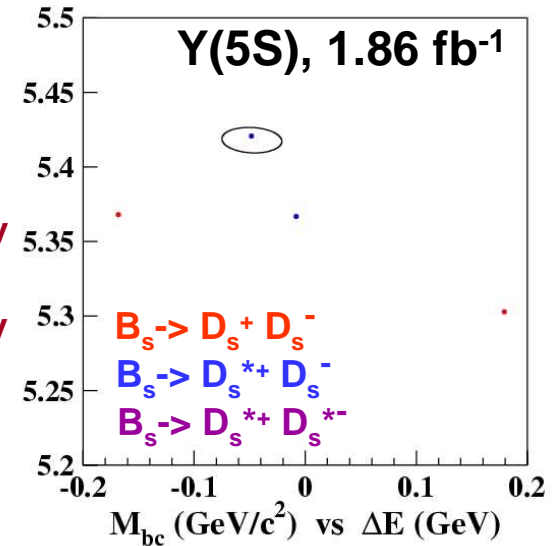
Expected with 25 fb⁻¹ at Y(5S):

$$\text{Eff}(B_s \rightarrow D_s^+ D_s^-) \sim 2 \times 10^{-4} \quad N \sim 2.5 \times 10^6 \times 2 \times 10^{-4} \times 10^{-2} \sim 5 \text{ ev}$$

$$\text{Eff}(B_s \rightarrow D_s^{*+} D_s^-) \sim 1 \times 10^{-4} \quad N \sim 2.5 \times 10^7 \times 10^{-4} \times 2 \times 10^{-2} \sim 5+5 \text{ ev}$$

$$\text{Eff}(B_s \rightarrow D_s^{*+} D_s^{*-}) \sim 5 \times 10^{-5} \quad N \sim 2.5 \times 10^7 \times 5 \times 10^{-5} \times 3 \times 10^{-2} \sim 4 \text{ ev}$$

=> Accuracy of $Bf(B_s \rightarrow D_s^{(*)+} D_s^{(*)-})$ has to be ~30%.



$D_s^+ \rightarrow \phi\pi^+, K^{*0} K^+, K_s K^+$

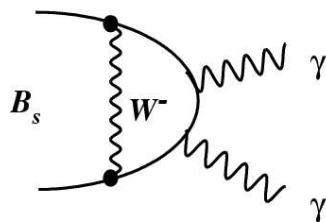
$$\frac{\Delta\Gamma_{CP^S}}{\Gamma_s} \approx \frac{Bf(B_s \rightarrow D_s^{(*)+} D_s^{(*)-})}{1 - Bf(B_s \rightarrow D_s^{(*)+} D_s^{(*)-}) / 2}$$

\Leftarrow should be compared with direct $\Delta\Gamma_s/\Gamma_s$ measurement to test SM.

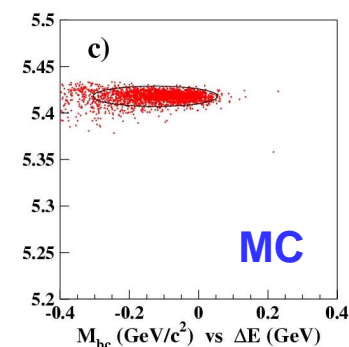
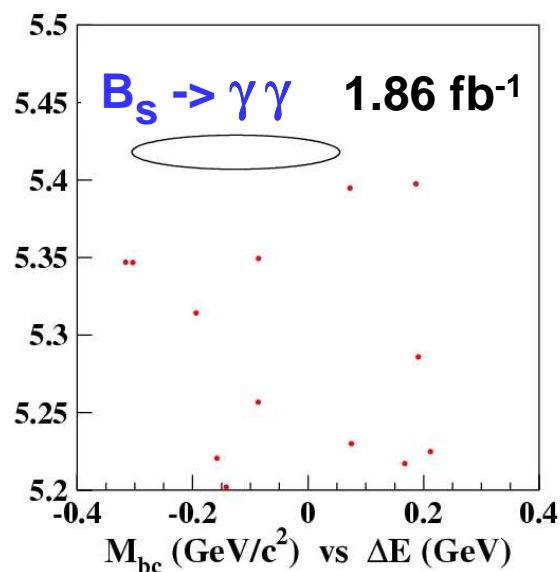
$\Delta\Gamma_s/\Gamma_s$ lifetime difference can be measured directly with high accuracy at Y(5S) and also at Tevatron and LHC experiments.



Exclusive $B_s \rightarrow \gamma\gamma$ decay



Natural mode to search for BSM effects, many theoretical papers devoted to this decay.



PDG limit : $Bf(B_s \rightarrow \gamma\gamma) < 1.48 \times 10^{-4}$

90% CL UL with 1.86 fb^{-1} : $Bf(B_s \rightarrow \gamma\gamma) < 0.53 \times 10^{-4}$.

Expected UL with 100 fb^{-1} : $Bf(B_s \rightarrow \gamma\gamma) < 1. \times 10^{-6}$.

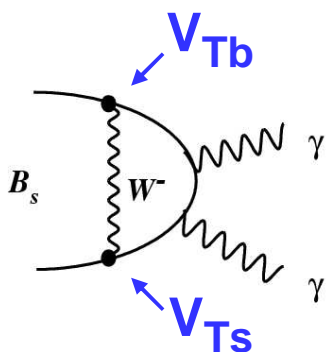
SM : $Bf(B_s \rightarrow \gamma\gamma) = (0.5-1.0) \times 10^{-6}$.

BSM can increase Bf up to two orders of magnitude.

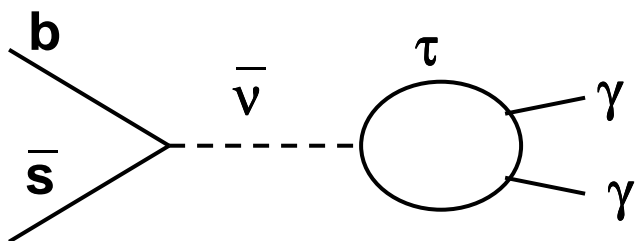


Exclusive $B_s \rightarrow \gamma\gamma$ decay

Many “conventional” BSM models can be better constrained by $B \rightarrow K^* \gamma$ and $B \rightarrow s \gamma\gamma$ processes, however not all. In some BSM models $B_s \rightarrow \gamma\gamma$ provides the best limit, in particular in 4-generation model (V_{Ts}) and R-parity violating SUSY ($\times M(\gamma\gamma)$).



hep-ph/0302177 (Huang et al) : limits on four-generation matrix elements V_{Tb} and V_{Ts} were obtained from B decays. $Bf(B_s \rightarrow \gamma\gamma)$ can increase up to one order of magnitude under specific conditions. Decay $B \rightarrow \gamma\gamma$ is not affected.



hep-ph/0404152 (Gernintern et al): within R-parity violation SUSY, diagram with sneutrino will increase $Bf(B_s \rightarrow \gamma\gamma)$ up to one order of magnitude.



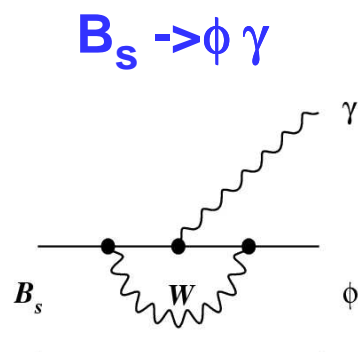
Exclusive $B_s \rightarrow K^- \pi^+$ and $B_s \rightarrow \phi \gamma$ decays

Direct CP-violation \Rightarrow non-zero charge asymmetry parameter
 $A = (Bf(+)-Bf(-)) / (Bf(+)+Bf(-))$.

Expected number of events with 5 ab^{-1} :

$B_s \rightarrow K^- \pi^+ : \sim 300$ events.

Large statistics is required to observe direct CP-violation on the level of 10%.



New Physics can contribute in penguin decay loops.

Estimated number of events with $100 \text{ fb}^{-1} : \sim 20$ ev.

Partner of $B \rightarrow K^* \gamma$ penguin decay .
Bf's have to be compared.



What else can be done at Super B Factory?

PDG ($Z \rightarrow b\bar{b}$, pp at $S^{1/2} = 1.8 \text{ TeV}$)

b hadron	fraction(%)
B^+ , B^0	39.8 ± 1.0
B_s	10.4 ± 1.4
b baryons	9.9 ± 1.7

Rates at e^+e^- continuum should be similar, baryon production is large.

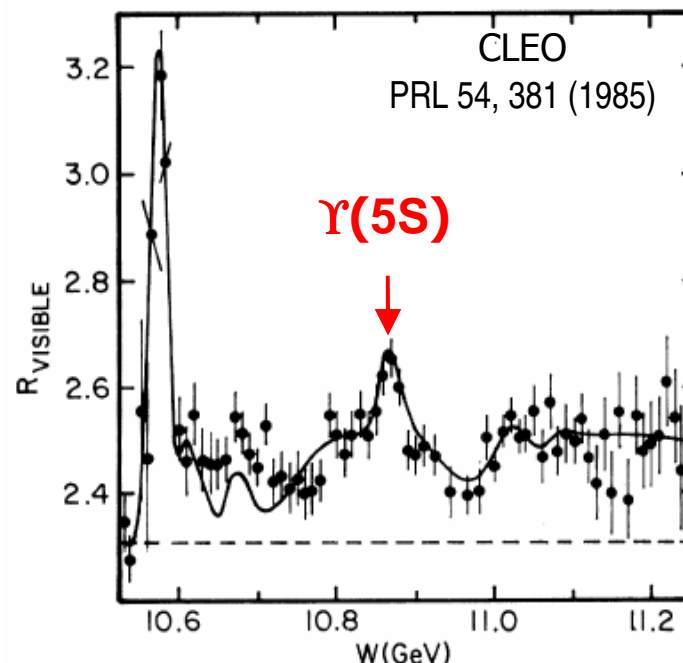
$$M(\Lambda_b) = (5624 \pm 9) \text{ MeV}/c^2$$

$$M(\Lambda_b) \times 2 = (11248 \pm 18) \text{ MeV}/c^2 \Rightarrow 6.3 \% \text{ up from } Y(4S) \text{ CME.}$$

Can Super B factory CM energy range be increased ?

$$M(B_c) = (6286 \pm 5) \text{ MeV}/c^2$$

$$e^+e^- \rightarrow B_s \bar{B}_s, \Lambda_b \bar{\Lambda}_b, B_c \bar{B}_c, \Xi_b \bar{\Xi}_b \dots ?$$





Conclusions

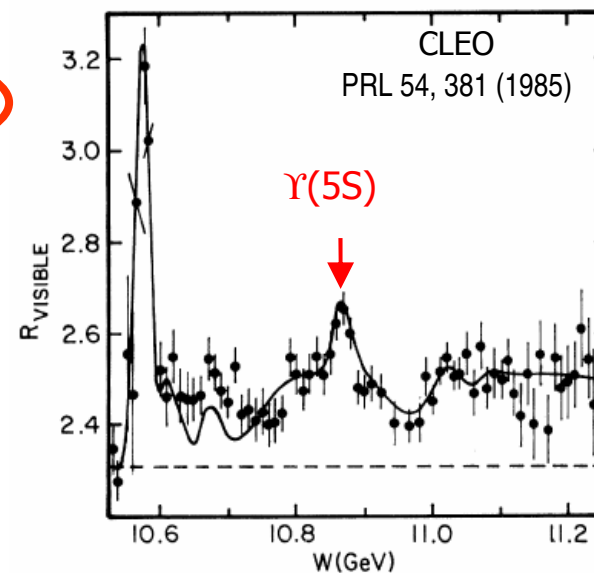
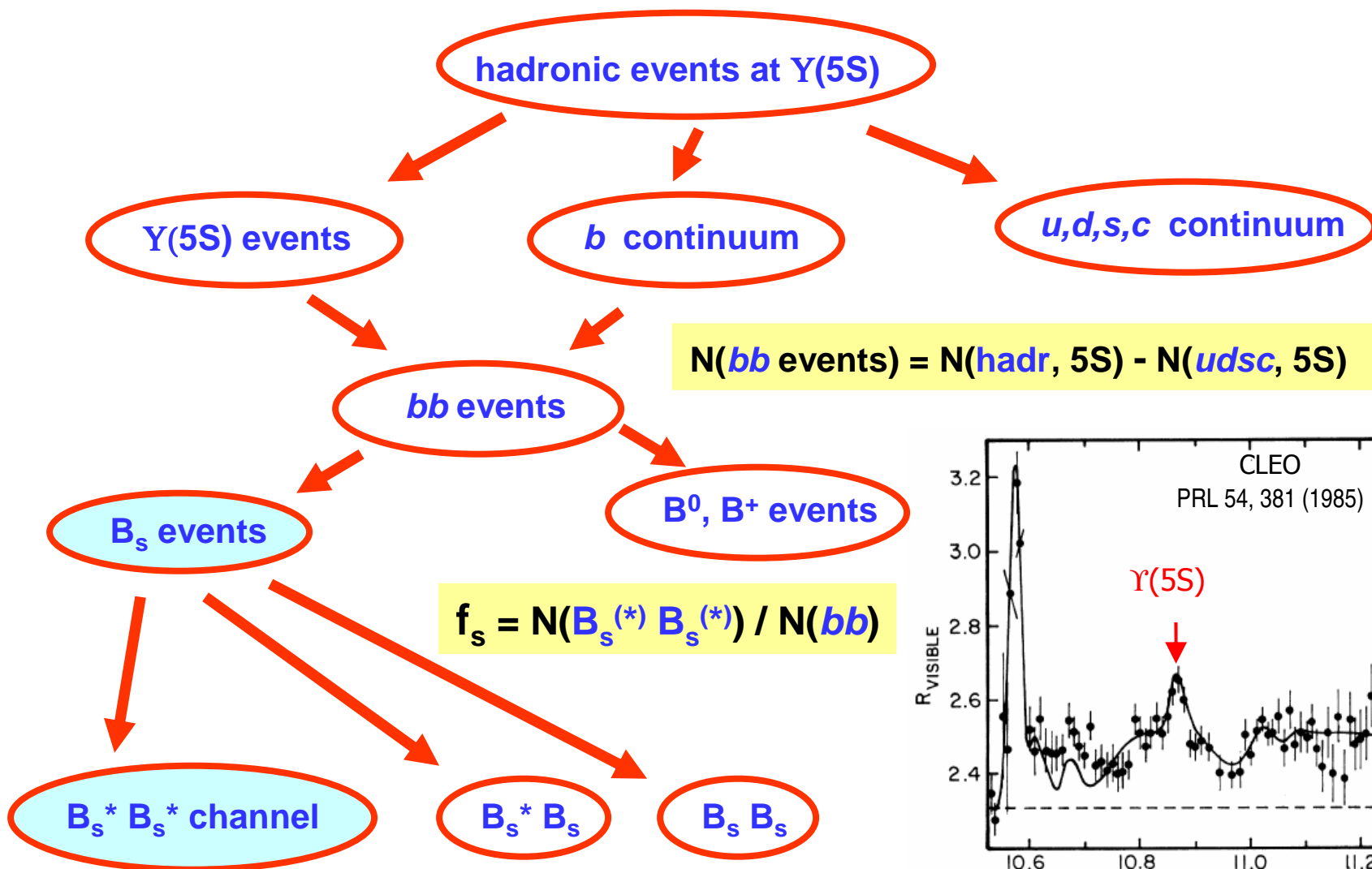
- B_s decays with branching fractions down to 10^{-6} can be studied with statistics of $\sim 100 \text{ fb}^{-1}$ at e+e- colliders running at $\Upsilon(5S)$.
- B_s studies at e+ e- colliders running at $\Upsilon(5S)$ have many advantages comparing with hadron-hadron colliders: high efficiency of photon reconstruction; 100% trigger efficiency; good K/π PID, partial reco.
- Many important SM tests can be done with statistics of the order of $(100-500) \text{ fb}^{-1}$ at Super B factory.
- Possibility to increase CM energy range (up to $\sim 14 \text{ GeV}$) should be considered for Super B factory. It could provide opportunity to study wide spectrum of b-hadrons, like Λ_b , B_c , Ξ_b .



Background slides



Hadronic event classification





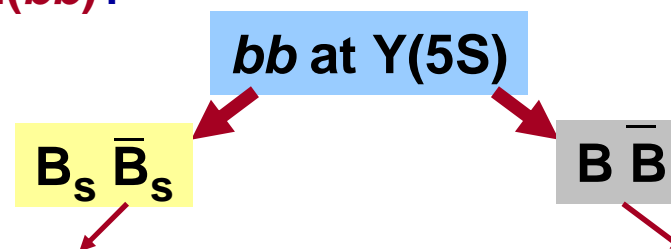
Number of bb events, number of B_s events

Y(5S) : Lumi = 1.857 ± 0.001 (stat) fb^{-1}

Cont (below 4S) : 3.670 ± 0.001 (stat) fb^{-1}

$N_{bb}(5S) = 561,000 \pm 3,000 \pm 29,000$ events \Rightarrow 5% uncertainty (from luminosity ratio)

How to determine $f_s = N(B_s^{(*)} B_s^{(*)}) / N(bb)$?



$$Bf(Y(5S) \rightarrow D_s X) / 2 = f_s \times Bf(B_s \rightarrow D_s X) + (1 - f_s) \times Bf(B \rightarrow D_s X)$$

1. $Bf(B_s \rightarrow D_s X)$ can be predicted theoretically, tree diagrams, large.
2. $Bf(B \rightarrow D_s X)$ is well measured at the Y(4S).

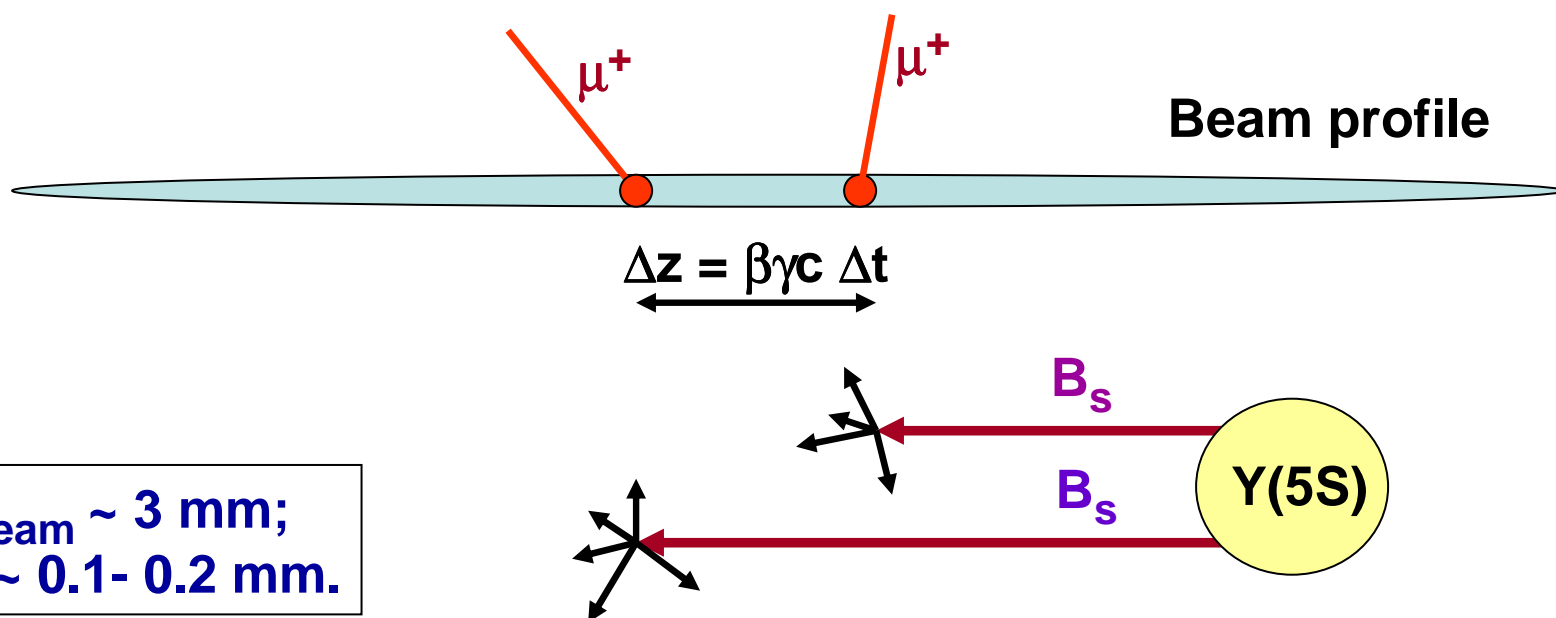


Feasibility of B_s lifetime measurement with same sign leptons

Lifetime can be measured using two fast same sign lepton tracks and beam profile. To remove secondary D meson semileptonic decays: $P(\ell) > 1.4 \text{ GeV}$.

$$Y(5S) : B_s(\ell^+) B_s(\ell^+) / B_s(\ell^+) B_s(\ell^-) = 100\%$$

$$Y(5S) : B(\ell^+) B(\ell^+) / B(\ell^+) B(\ell^-) \sim 10\%$$





Comparison with Fermilab B_s studies.

- There are several topics, where Y(5) running has advantages comparing with CDF and D0:
 - 1) **Model independent** branching fraction measurements.
 - 2) Measurement of decay modes with γ , π^0 and η in final state ($D_s^+ \rho^-$).
 - 3) Measurement of **multiparticle** final states (like $D_s^+ D_s^-$).
 - 4) **Inclusive** branching fraction measurements (semileptonic B_s).
 - 5) Partial reconstruction ($B_f(D_s^+ l^- \nu)$) using “missing-mass” method).

- There are also disadvantages:
 - 1) We have to choose between running at **Y(4S)** or **Y(5S)**.
 - 2) **Number of B_s** is smaller than in Fermilab experiments.
 - 3) Vertex resolution is **not** good enough to measure **B_s mixing**.

- **Running at Y(5S) is a new tool to study B_s physics. First Belle results are very promising. Collider exists => relatively low price.**